Ultra High-Precision Astrometry with the Habitable Worlds Observatory

Scott Gaudi

The Ohio State University

With contributions from Eduardo Bendek, Kaz Gary (OSU Graduate Student), Aki Roberge, Breann Starski

A Future Space Mission with Very High Precision Astrometry Workshop

Institut d'Astrophysique de Paris

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PRELIMINARY SPECS & CANDIDATE INSTRUMENTS

UV Multi-Object Spectrograph UV/Vis multi-object spectroscopy and FUV imaging Bandpass \vert ~100–1000 nm Field-of-View \sim 2' \times 2' Apertures $\sim 840 \times 420$ $R(\lambda/\Delta\lambda)$ ~500–60,000

TECHNICAL ACTIVITIES – EXPLORATORY ANALYTIC CASES (EACS)

These are 1st round mission architectures that will be used to explore the HWO trade space. Purposes ...

- Practice end-to-end modeling, from science to engineering. Develop initial codes to "pipeclean" the process
- Use EACs to identify key technology gaps and guide maturation of potential technology solutions
- Provide feedback to rocket vendors as soon as possible to help influence their direction

Exploration of three Round 1 EACs will take \sim 1 year. Findings will fold into Round 2 EACs.

6-m inner diameter / 7.2-m outer diameter off-axis

8-m diameter on-axis

HIGH RESOLUTION IMAGER

EAC1 HRI NIR INSTRUMENT LAYOUT

Masses of Nearby Habitable Planets: Astrometry, Direct Imaging and RV

Courtesy of Eduardo Bendek

Courtesy of Eduardo Bendek

Astrometry, Direct Imaging and RV

Plot by Eric Mamajek

Astrometric Detection of Exoplanets

- This approximate $\frac{s}{N}$ assumes *N* uniformsampled 2D astrometric measurements each with precision σ along each axis,.
- The astrometric signal of planets with periods longer than the duration of the survey T are very hard to detect
- The signals are mostly linear 'trends', which cannot be distinguished from the host star proper motion.

$$
S_{\bigwedge} = \begin{cases} \sqrt{N}F(i, e, \omega) \frac{\alpha_0}{\sigma} & \text{if } P \le T \\ 0 & \text{if } P \ge T \end{cases}
$$

where $\alpha_0 \equiv \left(\frac{a}{d}\right)$ \boldsymbol{d} m_p M_* is the astrometric signal for a circular, face-on orbit, and

$$
F(i, e, \omega) = \sqrt{\frac{1}{2} (1 + \cos^2 i) \left[1 - e^2 \left[3 - \left(\frac{2}{1 + \cos^2 i} - 1 \right) \right] \cos^2 \omega \right]}
$$

accounts for the effects of inclination i , eccentricity e , and longitude of periastron ω , and is generally of order unity.

HZ planets: signal and comparison to RV

- The signal for an Earth analog at 10 pc is 0.3 μ as.
- For HZ planets, assuming $L_* \propto M_*^4$, the signal increases as M_*^2 .
- All else fixed, planets in the HZ are easier to detect around higher mass stars.
- For RV, the signal goes as $M_*^{-3/2}$, so HZ planets are easier to detect around lower-mass stars.
- High mass stars also have fewer spectral lines and ten to be rapidly rotating (above the Kraft break).

$$
\alpha_0 \equiv 0.3 \mu as \left(\frac{a}{AU}\right) \left(\frac{m_p}{M_{\oplus}}\right) \left(\frac{M_*}{M_{\odot}}\right)^{-1} \left(\frac{d}{10 pc}\right)^{-1}
$$

$$
\alpha_0 \sim 0.3 \mu \text{as} \left(\frac{m_p}{M_{\oplus}}\right) \left(\frac{M_*}{M_{\odot}}\right) \left(\frac{d}{10 pc}\right)^{-1} \text{for HZ}
$$

$$
K \sim 10 \text{ cm/s} \left(\frac{m_p}{M_{\oplus}}\right) \left(\frac{M_*}{M_{\odot}}\right)^{-3/2} \text{for HZ}
$$

11

Detection and Mass Uncertainty

$$
S_{\text{N}} \simeq 10F \left(\frac{N}{100}\right)^{1/2} \left(\frac{\sigma}{0.3 \mu \text{as}}\right)^{-1} \left(\frac{m_p}{M_{\oplus}}\right) \left(\frac{M_*}{M_{\odot}}\right)^2 \text{ if } P \lesssim T
$$

- A planet $m_p = M_{\oplus}$ planet in the HZ can be detected with $^S\text{/}_N \simeq 10$ with N=100 (2D) astrometric measurements each with precision $\sigma = 0.3 \ \mu$ as along each axis.
- This corresponds to a 10% mass measurement if the observations are spread out uniformly in phase and the duration of observation is (much) greater than the period.
- Highly simplified treatment but provides a rough sense of the number and quality of the astrometric measurements needed.

Photon Noise Astrometric Uncertainty

$$
\sigma_a \simeq 1.7 \mu \text{as} \ 10^{-0.2(V_{ref}-15)} \left(\frac{t_{exp}}{300 \text{s}}\right)^{-1/2} \left(\frac{D}{6.5 m}\right)^{-1} \left(\frac{\lambda/D}{0.0159 \text{ as}}\right)
$$

- The astrometric uncertainty for a V_{ref} = 15 reference star is \sim 1.7 μ as for a 5-minute exposure on HWO in the V band.
- Need ~30 reference stars with $V_{ref} \le 15$ to achieve the total astrometric uncertainty of \sim 0.3 μ as per epoch.

Number of reference stars

- For HWO, the number of reference stars in the FOV of the camera is a critical consideration
- Depends on Galactic latitude and longitude and filter
- Generally, follows a broken power law.

$$
\Sigma(V) = \Sigma_b \left[10^{-s\alpha_1(V-V_b)} + 10^{-s\alpha_2(V-V_b)} \right]^{-1/s}
$$

- α_1 is the bright-end slope, α_2 is the faint-end slope, Σ_b is the # stars/deg²/mag at $V = V_b$, and *s* is the 'softness' of the break
- α_1 ~0.4 implies a uniform number of density of stars of a fixed magnitude
- Broken power law due to disk geometry.

Jumber

Number of reference stars

• Can use this form to estimate the effective number of reference stars in any given line-of-sight.

$$
\Sigma(V) = \Sigma_b \left[10^{-s\alpha_1(V-V_b)} + 10^{-s\alpha_2(V-V_b)} \right]^{-1/s}
$$

- If $\alpha_1 = 0.4$, $\alpha_1 \sim 0$, and $s \gg 1$, the there is an equal contribution per mag for $V < V_b$, and the effective number of reference stars is $N_{\text{ref.eff}} = \Sigma_b V_b$
- For b~20°, α_1 ~0.4, α_2 ~0.0, V_b ~15, Σ_b ~10³deg⁻²mag⁻¹, so $N_{\rm ref,eff}$ ~ 10⁴ deg⁻²
- Number of effective reference stars will be much lower at higher Galactic latitudes
- Likely need a FOV of \sim 10' \times 10' to do Earths around the high-latitude starget stars
- \sim 2 \times 5 times larger than nominal imaging camera.
- Requires roughly 64 8k by 8k CMOS detectors for Nyquist sampling (!)

Scaling with size of the detector

$$
S_{N} \sim f \alpha \Gamma_{\gamma,0}^{1/2} N_{\rm ref,eff}^{1/2} 10^{-0.2 V_b} \ell_d t_{exp}^{1/2} N_d^{1/2} f_{\#}^{-1} (\lambda/D)^{-1}
$$

- α is the magnitude of the astrometric signal
- $\Gamma_{\nu,0}$ is the photon rate per area per time a V=0
- $N_{\text{ref.eff}}$ is the effective number of reference stars
- V_b is the break magnitude
- $t_{exp}^{1/2}$ is the exposure time
- $N_d^{1/2}$ is the number of data points
- \bullet *D* is the diameter
- λ is the wavelength
- ℓ_d is the linear size of the detector = $(Area)^{1/2}$
- $f_{\#}$ is the f-number of the system
- 1011 • f is a dimensionless constant of order unity, that depends on the Keplerian parameters, shape of the PSF, relation between the diameter and effective area of the telescope

Scaling with size of the detector

$$
S_{N} \sim f \alpha \Gamma_{\gamma,0}^{1/2} N_{\rm ref,eff}^{1/2} 10^{-0.2 V_b} N_{pix}^{1/2} t_{exp}^{1/2} N_d^{1/2} D
$$

- α is the magnitude of the astrometric signal
- $\Gamma_{\nu,0}$ is the photon rate per area per time a V=0
- $N_{\text{ref.eff}}$ is the effective number of reference stars
- V_b is the break magnitude
- $t_{exp}^{1/2}$ is the exposure time
- $N_d^{1/2}$ is the number of data points
- \bullet *D* is the diameter
- N_{pix} is the number of pixels
- \bullet f is a dimensionless constant of order unity that depends on the Keplerian parameters, relation between the effective area of the telescope and the diameter.
- The wavelength cancels out: smaller wavelength→smaller PSF → smaller pixels for Nyquist sampling → small FOV for fixed # of pixels

Astrometric Microlensing

Background stars will be deflected by an amount:

$$
\theta_E \approx \sqrt{\kappa M \pi} \sim 0.03'' \left(\frac{M}{M_{\odot}}\right)^{1/2} \left(\frac{\pi}{0.1''}\right)^{1/2}
$$

• where $\kappa=8.14$ mas M_{\odot}^{-1}

 $\delta(\theta) \sim \frac{\theta_E^2}{\theta}$

 θ

- The typical angular distance to a reference star is $\sim \theta_d/\sqrt{2}$, where $\theta_d \sim (FOV)^{1/2} \sim 150''$
- The typical size of the astrometric shift is \sim 8 μ as.
- Thus, the S/N for all reference stars and all measurements is

$$
S_{\text{N}} \simeq 800 \left(\frac{M}{M_{\odot}}\right)^{1/2} \left(\frac{\pi}{0.1}\right)^{1/2} \left(\frac{N_e}{100}\right)^{1/2} \left(\frac{N_{\text{eff,ref}}}{100}\right)^{1/2} \left(\frac{\text{FOV}}{100 \text{ sq.armin}}\right)^{1/2}
$$
\n• The uncertainty in the mass is $\frac{\delta M}{M} \sim (S/N)^{-1} \sim 0.1\%$

Background due to halo of the target star

• At 1', the flux from the halo of the target star in the PSF of the reference stars is 27 magnitudes fainter!

EAC1 HRI UVIS DISTORTION

Questions

- How well do we need to measure the masses of the planets? (Note: We will also likely get a measurement of the mass of the star.)
- How well we know the orbits? Can this reduce the amount of time required to get masses?
- What wavelengths are best? Are broader or narrower filters better?
- Can we use the guide camera?
- How do we deal with the dynamic range issue?
- How can we control systematics at the \sim 2×10⁻⁶ of a pixel level?